LTC Fresnel Approximation Stephen Hill

To explain our Fresnel approximation, I will start first of all with a simpler case that we will then build upon. Let's imagine that we have a microfacet BRDF, $\rho(\boldsymbol{\omega}_{\boldsymbol{v}}, \boldsymbol{\omega}_{\boldsymbol{l}})$, with shadowing-masking but without Fresnel. Due to the presence of shadowing, the integral of the cosine-weighted BRDF over the sphere can be less than 1 (i.e. some 'energy' has been lost due to shadowing):

$$\int_{\Omega} \rho(\boldsymbol{\omega}_{\boldsymbol{v}}, \boldsymbol{\omega}_{\boldsymbol{l}}) \cos \theta_{\boldsymbol{l}} \, \mathrm{d} \boldsymbol{\omega}_{\boldsymbol{l}} \le 1.$$
(1)

In contrast and by design, LTCs always integrate to 1. Therefore, in order to achieve an accurate fit of the BRDF using LTCs, we store the norm¹ n_D (the magnitude of the BRDF) as a scale factor, in addition to M^{-1} (which captures the shape of the BRDF):

$$n_D = \int_{\Omega} \rho(\boldsymbol{\omega}_{\boldsymbol{v}}, \boldsymbol{\omega}_{\boldsymbol{l}}) \cos \theta_l \, \mathrm{d}\boldsymbol{\omega}_{\boldsymbol{l}}.$$
⁽²⁾

As with M^{-1} , we store n_D in a 2D texture, parameterized by incident direction and roughness.

Now let's turn our attention to Fresnel. Rather than attempting to fit LTCs to a BRDF that directly incorporates Fresnel, we instead choose to treat the Fresnel term separately, as an additional influence on the magnitude of the BRDF. Thus the norm becomes

$$\int_{\Omega} F(\boldsymbol{\omega}_{\boldsymbol{v}}, \boldsymbol{\omega}_{\boldsymbol{l}}) \,\rho(\boldsymbol{\omega}_{\boldsymbol{v}}, \boldsymbol{\omega}_{\boldsymbol{l}}) \cos \theta_l \,\mathrm{d}\boldsymbol{\omega}_{\boldsymbol{l}}. \tag{3}$$

Using Schlick's approximation [Sch94], this expands to

$$\int_{\Omega} [R_0 + (1 - R_0)(1 - \langle \boldsymbol{\omega}_{\boldsymbol{v}}, \boldsymbol{\omega}_{\boldsymbol{h}} \rangle)^5] \rho(\boldsymbol{\omega}_{\boldsymbol{v}}, \boldsymbol{\omega}_{\boldsymbol{l}}) \cos \theta_l \, \mathrm{d}\boldsymbol{\omega}_{\boldsymbol{l}}, \tag{4}$$

which we can rearrange to the following:

$$= R_0 \int_{\Omega} \rho(\boldsymbol{\omega}_{\boldsymbol{v}}, \boldsymbol{\omega}_{\boldsymbol{l}}) \cos \theta_l \, \mathrm{d}\boldsymbol{\omega}_{\boldsymbol{l}} + (1 - R_0) \int_{\Omega} (1 - \langle \boldsymbol{\omega}_{\boldsymbol{v}}, \boldsymbol{\omega}_{\boldsymbol{h}} \rangle)^5 \rho(\boldsymbol{\omega}_{\boldsymbol{v}}, \boldsymbol{\omega}_{\boldsymbol{l}}) \cos \theta_l \, \mathrm{d}\boldsymbol{\omega}_{\boldsymbol{l}}$$
$$= R_0 \, n_D + (1 - R_0) \, f_D, \quad \text{where } f_D = \int_{\Omega} (1 - \langle \boldsymbol{\omega}_{\boldsymbol{v}}, \boldsymbol{\omega}_{\boldsymbol{h}} \rangle)^5 \rho(\boldsymbol{\omega}_{\boldsymbol{v}}, \boldsymbol{\omega}_{\boldsymbol{l}}) \cos \theta_l \, \mathrm{d}\boldsymbol{\omega}_{\boldsymbol{l}}. \tag{5}$$

Now, in addition to n_D that we had before, we store a second term f_D^2 .

While this may seem like a coarse approximation, we found it to work very well in practice from a visual standpoint, and it's an approach that has already proven to be effective in the context of environmental illumination [Kar13; Laz13]. Furthermore, this solution avoids complicating the fitting process or increasing the dimensionality of the tabulated data.

¹This is mentioned very briefly in our paper at the start of the Representation and Storage section: "Furthermore, we use an additional parameter for the norm \dots ".

²We could alternatively store $n'_D = n_D - f_D$, and calculate $R_0 n'_D + f_D$ at runtime, which saves an ALU instruction.

References

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